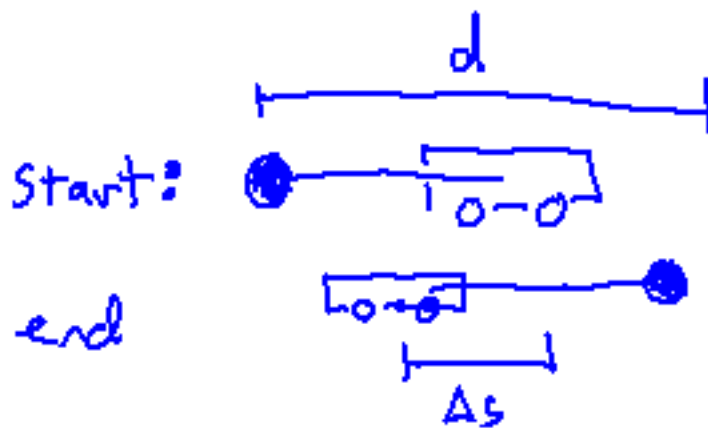


Momentum Matching for the Pendulum on Cart Problem

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1 Problem Description



motion



There is a pendulum with mass m_1 that is connected with a cart of mass m_2 . The object is to move the cart a distance Δs ,



when the pendulum position moves from $-\frac{\pi}{2}$ to $\frac{\pi}{2}$. The coordinate system that will be used is:

1.1 Problem Setup

We first need to set up the equations of motion for this problem. So we need to set up the lagrangian of the system. The lagrangian of a general system is: $KE - U = L$

And for our system:

$$L = \frac{1}{2}\dot{\theta}^2 l^2 m_2 + s\dot{\theta}l \cos \theta m_2 + \frac{1}{2}\dot{s}^2 (m_1 + m_2) - m_2 l g \cos(\theta). \quad (1)$$

In order to acquire the equations of motion:

$$\frac{\partial}{\partial t} \left(\frac{\partial L}{\partial \dot{\theta}} \right) - \frac{\partial L}{\partial \theta} = 0$$

$$\frac{\partial}{\partial t} \frac{\partial L}{\partial \dot{s}} = 0$$

2 Controller

In order to meet our goal, there are two possible controls that could be possible: controlling the pendulum arm or controlling the cart.

2.1 Pendulum Arm Controller

If we were to control the pendulum arm:

$$\ddot{\theta} l^2 m_2 + \dot{s} \cos \theta m_2 + g \sin \theta m_2 = u \quad (2)$$

$$\ddot{\theta} l^2 \sin \theta m_2 + \dot{s} (m_1 + m_2) = 0 \quad (3)$$

However due to the conservation of linear momentum, no matter what controls are applied to the pendulum arm or whatever path θ makes, if $\theta(T)$ ends up at $\frac{\pi}{2}$ at some point T , the displacement of s is fixed. Assuming that the system begins from rest:

$$\dot{s} = \frac{-lm_2}{m_1 + m_2} \cos(\theta) \dot{\theta} = \frac{-lm_2}{m_1 + m_2} \frac{d}{dt} \sin(\theta(t)) \quad (4)$$

So if $\theta(0) = -\frac{\pi}{2}$, then the $\Delta s = -2 \frac{-lm_2}{m_1 + m_2}$.

2.2 Controlling the Cart

Therefore, our only option is to control the cart.

$$\ddot{\theta} l^2 m_2 + \dot{s} l \cos(\theta) m_2 + g l \sin(\theta) m_2 = 0 \quad (5)$$

$$\ddot{\theta} \cos(\theta) m_2 - d\theta^2 l \sin(\theta) m_2 + \dot{s} (m_1 + m_2) = u \quad (6)$$

2.3 Momentum Matching

Based on the behavior of the uncontrolled system, we know that the system displaces by $-2 \frac{m_2 l}{m_1 + m_2}$, and if we substitute $\beta = m_2 l$ and $\gamma = m_1 + m_2$, the uncontrolled system displaces by $\frac{\beta}{\gamma}$. This is because of the conservation of momentum equation. So, to consider whether or not a control exists, so that the system conserves a different linear momentum equation, we need to find a control u such that the momentum is conserved:

$$\dot{\theta} \beta_n \cos(\theta) + \dot{s} \gamma_n = C$$

. with C being a constant.

This has a time derivative of:

$$\ddot{\theta} \cos(\theta) \beta_n - \dot{\theta}^2 \sin(\theta) \beta_n + \dot{s} \gamma_n = 0 \quad (7)$$

Using the controlled equations above:

$$\ddot{\theta} + \frac{\dot{\theta}^2 \cos(\theta) \sin(\theta) m_2}{m_1 + \sin(\theta)^2 m_2} + \frac{g \sin(\theta) \gamma}{m_1 l + \sin(\theta)^2 \beta} = \frac{u \cos(\theta)}{l m_1 + l \sin(\theta)^2} \quad (8)$$

$$\ddot{s} - \frac{\dot{\theta}^2 l \sin(\theta) m_2}{m_1 + \sin(\theta)^2 m_2} - \frac{g \cos(\theta) \sin(\theta) m_2}{m_1 + \sin(\theta)^2 m_2} = \frac{u}{m_1 + \sin(\theta)^2 m_2} \quad (9)$$

To replace the accelerations one sees a control of the form

$$u = u_0(\theta) + u_1 \dot{\theta}^2 \quad (10)$$

Solving for u_0 and $u_1(\theta)$ will result in:

$$u_1 = \frac{l \sin(\theta) (\kappa m_1 + (-l + \kappa) m_2)}{l - \kappa \cos(\theta)^2} \quad (11)$$

$$u_0 = \frac{g \sin(2\theta) (l m_2 - \kappa (m_1 m_2))}{-2l + \kappa + \kappa \cos(2\theta)} \quad (12)$$

For $\kappa < 1$, the denominators are nonzero and the control is bounded. Using this controller, the equations become:

$$\ddot{\theta} = \frac{2(g + d\dot{\theta}^2 \kappa \cos(\theta) \sin(\theta))}{-2l + \kappa + \kappa \cos(2\theta)} \quad (13)$$

$$\ddot{s} = -\frac{2\kappa(d\dot{\theta}^2 l + g \cos(\theta)) \sin(\theta)}{-2l + \kappa + \kappa \cos(2\theta)} \quad (14)$$

If the pendulum arm goes from $\frac{-\pi}{2}$ to $\frac{\pi}{2}$, then the cart will displace a distance of 2κ to $\frac{\pi}{2}$, or in other words, if the initial conditions are as such: $\theta(0) = \frac{-\pi}{2}$; $\dot{\theta}(0) = 0$; $s(0) = 0$; there must be a time τ , when $\theta(\tau) = \frac{\pi}{2}$; $\dot{\theta} = \frac{\pi}{2}$; $s = 0$; So now, we will use a concept called time-reversal symmetry and symmetry over θ , where $\theta = -\theta$. The concept of time-reversal symmetry means that a system will follow the same path regardless of whether time is positive or negative. Using this, we will show that if the pendulum moves from $\frac{-\pi}{2}$ to 0. Then it implies that the pendulum can also move from 0 to $\frac{-\pi}{2}$. The second symmetry will imply that if the pendulum moves from $\frac{-\pi}{2}$ to 0 then it will also move from $\frac{\pi}{2}$ to 0, thus it with these two symmetries combined, we will show that the pendulum will move from $\frac{-\pi}{2}$ to $\frac{\pi}{2}$

Proof 1 : For a system to have time-reversal symmetry, there must exist a time T when $\theta(T) = 0$, and it is invariant under the transformations $t \rightarrow -t$, and $v \rightarrow -v$. Our system is:

$$\dot{\theta} = V \quad (15)$$

$$\dot{V} = \frac{2(g + V^2 \kappa \cos(\theta)) \sin(\theta)}{-2l + \kappa + \kappa \cos(2\theta)} \quad (16)$$

a) If we make a change from $t \rightarrow -t$ then we must show that if $(\theta(t), V(t))$ is a solution, then so must $(\theta(-t), -V(-t))$. So then

$$\frac{\partial}{\partial t}(\theta(-t), -V(-t)) = (-\dot{\theta}(-t), \dot{V}(-t))$$

Now, if substitute these values into the equation above:

$$-\dot{\theta}(-t) = -V(-t) \quad (17)$$

$$\dot{V}(-t) = \frac{2(g + (-V(-t))^2 \kappa \cos(\theta(-t)) \sin(\theta(-t)))}{-2l + \kappa + \kappa \cos(2\theta(-t))} \quad (18)$$

Because of the nature of the pendulum that can be seen by the above equation, $\theta(t) = \theta(-t)$ and also in the equation above, $(-V(-t))$ is squared. Since the $y = x^2$ is an even function, The equation becomes:

$$-\dot{\theta}(-t) = -V(-t) \quad (19)$$

$$\dot{V}(-t) = \frac{2(g + (V(-t))^2 \kappa \cos(\theta(t)) \sin(\theta(t)))}{-2l + \kappa + \kappa \cos(2\theta(t))} \quad (20)$$

So with the time reversal symmetry, if $\theta(\tau) = 0$ for some time τ , when reversed, the system will be invariant under the transformation $t \rightarrow -t$. It will simply move in the opposite direction, while the acceleration remains invariant.

b) If the system were to be transformed from $\theta(t) = -\theta(t)$, then we can deduce that if the pendulum takes the path from $-\frac{\pi}{2}$ to 0. That it will take the path from 0 to $\frac{\pi}{2}$ then $-\dot{\theta} = -V(t)$. If we replace these in the equation above we get:

$$-\dot{\theta} = -V \quad (21)$$

$$\dot{V}(t) = \frac{2(g + (-V)^2 \kappa \cos(\theta(t)) \sin(\theta(t)))}{-2l + \kappa + \kappa \cos(2\theta(t))} \quad (22)$$

The acceleration remains the same because $V^2 = (-V)^2$. This shows that if θ goes from $-\frac{\pi}{2}$ to 0 then it will also go from 0 to $\frac{\pi}{2}$

c) Now, we must show that there is at time T such that $\theta(T) = 0$. If you look at the acceleration equations:

$$\ddot{\theta} = \frac{2(g + \dot{\theta}^2 \kappa \cos(\theta)) \sin(\theta)}{-2l + \kappa + \kappa \cos(2\theta)} \quad (23)$$

$$\ddot{s} = -\frac{2\kappa(d\theta^2 l + g \cos(\theta)) \sin(\theta)}{-2l + \kappa + \kappa \cos(2\theta)} \quad (24)$$

at $T = 0$ $\ddot{\theta} = 0$. Now, as T goes from $-\frac{\pi}{2}$ to 0, the acceleration continues to remain positive. So, by deduction, it can be seen that if a system starts at rest and it continues to have positive acceleration until $t = 0$, then it must pass through that point. Therefore by using time reversal symmetry, we have shown that momentum matching is a viable solution, when the pendulum goes from $-\pi/2$ to $\pi/2$.

3 Future Problems

We then tried to see if there was a way using this symmetry to get the pendulum to move from $-\pi/2$ to $\pi/2 + k$ although there is symmetry about any two points K and $-K$. The acceleration does not accelerate more than at point $K + -\frac{\pi}{2}$, therefore, we must try to match other functions onto u_1 and u_0 .